

Continuous symmetries

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- spacetime translations $x^M \rightarrow x'^M = x^M + \epsilon^M$
 $\Rightarrow \delta\psi = \epsilon^m \partial_m \psi$
 $\Rightarrow T^{mv} = i \bar{\psi} \gamma^m \partial^v \psi - g^{mv} \mathcal{L}$
 \uparrow the energy-momentum tensor

The equations of motion are already implied in the (derivation of the) conserved current, so we can choose $\mathcal{L} = 0$ and

$$T^{mv} = i \bar{\psi} \gamma^m \partial^v \psi$$

Then the total energy is

$$\begin{aligned} E &= \int d^3\bar{x} T^{00} \\ &= i \int d^3\bar{x} \bar{\psi} \partial^0 \psi \\ &= \int d^3\bar{x} \bar{\psi} (-i \gamma^i \partial_i + m) \psi \end{aligned}$$

- Lorentz transformations $x^M \rightarrow x'^M = \Lambda^M_{\nu} x^{\nu}$
 $\Rightarrow \delta\psi^{\alpha} = -\omega_{\mu\nu} (x^{\nu} \partial^{\mu} \psi^{\alpha} - \frac{i}{2} (\Sigma^{\mu\nu})^{\alpha}_{\beta} \psi^{\beta})$
 $\Rightarrow (J^{\mu\nu})^{\alpha\beta} = x^{\nu} T^{\mu\alpha} - x^{\mu} T^{\nu\alpha} + \bar{\psi} \gamma^{\mu} \Sigma^{\nu\alpha} \psi$

When we quantise, as we will see, $J^{\mu\nu}$ becomes an operator and this gives rise to an internal angular momentum (ie. a particle with spin $1/2$!)

• Vector symmetry

$$\Psi \rightarrow \Psi' = e^{-i\alpha} \Psi$$

$$\Rightarrow \delta \Psi = -i\delta\alpha \Psi$$

$$\Rightarrow j_\nu^\mu = \bar{\Psi} \gamma^\mu \Psi$$

N.B. $\partial_\mu j_\nu^\mu = (\partial_\mu \bar{\Psi}) \gamma^\mu \Psi + \bar{\Psi} \gamma^\mu \partial_\mu \Psi$
 $= im \bar{\Psi} \Psi - im \bar{\Psi} \Psi$
 $= 0$

↑ this is the conserved
vector current

The corresponding conserved charge is

$$Q = \int d^3\vec{x} \bar{\Psi} \gamma^0 \Psi = \int d^3\vec{x} \Psi^\dagger \Psi = \int d^3\vec{x} |\Psi|^2$$

↑
 electric charge
 (or, equivalently, particle #)

• Axial symmetry ($M=0$)

$$\Psi \rightarrow \Psi' = e^{i\alpha \gamma^5} \Psi, \bar{\Psi} \rightarrow \bar{\Psi}' = \bar{\Psi} e^{i\alpha \gamma^5}$$

N.B. $\partial_\mu j_A^\mu = (\partial_\mu \bar{\Psi}) \gamma^\mu \gamma^5 \Psi + \bar{\Psi} \gamma^\mu \gamma^5 \partial_\mu \Psi$
 $= 2im \bar{\Psi} \gamma^5 \Psi$

$$\Rightarrow \delta \Psi = i\delta\alpha \gamma^5 \Psi, \delta \bar{\Psi} = +i\delta\alpha \bar{\Psi} \gamma^5$$

$$\Rightarrow j_A^\mu = \bar{\Psi} \gamma^\mu \gamma^5 \Psi$$

↑
 note
 $e^{-i\alpha \gamma^5} \gamma^0 = \gamma^0 e^{i\alpha \gamma^5}$

This is one of the most interesting symmetries,
 use it is anomalous; it is a symmetry of
 the classical theory, but not of the quantum theory!

Here $\gamma^5 = i \gamma^0 \gamma^1 \gamma^2 \gamma^3$

$$(\gamma^5)^\dagger = \gamma^5$$

$$(\gamma^5)^2 = \mathbb{1}$$

$$\{\gamma^5, \gamma^\mu\} = 0$$

* Quick Quiz